Electromagnetic Resonators

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Waveguides (recap)



1. Maxwell:

(i)
$$\nabla \cdot \mathbf{E} = \mathbf{0}$$
, (iii) $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$,
(ii) $\nabla \cdot \mathbf{B} = \mathbf{0}$, (iv) $\nabla \times \mathbf{B} = \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}$.

2. Expand fields into components:

$$\vec{E}(x, y, z, t) = \vec{E}_0(x, y) e^{j(\omega t - kz)}$$
$$\vec{E}_0 = E_x \ \hat{x} + E_y \ \hat{y} + E_z \ \hat{z}$$

3. Transverse components in terms of Ez, Bz:

(i)
$$E_x = \frac{i}{(\omega/c)^2 - k^2} \left(k \frac{\partial E_z}{\partial x} + \omega \frac{\partial B_z}{\partial y} \right),$$

(ii) $E_y = \frac{i}{(\omega/c)^2 - k^2} \left(k \frac{\partial E_z}{\partial y} - \omega \frac{\partial B_z}{\partial x} \right),$
(iii) $B_x = \frac{i}{(\omega/c)^2 - k^2} \left(k \frac{\partial B_z}{\partial x} - \frac{\omega}{c^2} \frac{\partial E_z}{\partial y} \right),$
(iv) $B_y = \frac{i}{(\omega/c)^2 - k^2} \left(k \frac{\partial B_z}{\partial y} + \frac{\omega}{c^2} \frac{\partial E_z}{\partial x} \right).$

(Note: replace i by –j above)

4. Solve for Ez, Bz from here:

(i)
$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + (\omega/c)^2 - k^2\right] E_z = 0,$$

(ii)
$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + (\omega/c)^2 - k^2\right] B_z = 0.$$

- TE Solutions of the form
 - $B_{z}(x, y, z, t) = B_{0} \cos\left(\frac{m\pi x}{a}\right) \cos\left(\frac{n\pi y}{b}\right) e^{j(\omega t kz)}$ i.e. a fwd travelling wave



Cap the ends with metal plates [2]

$$B_{z} = \left(B_{0}^{+}e^{-jkz} + B_{0}^{-}e^{jkz}\right)\cos\left(\frac{m\pi x}{a}\right)\cos\left(\frac{n\pi y}{b}\right)e^{j\omega t}$$

Fwd and bkwd waves



Boundary conditions at z = 0, z = d: $(\nabla \times \vec{E})_z = -j\omega B_z$ $= \frac{\partial E_x}{\partial y} - \frac{\partial E_y}{\partial x}$ = 0

Cap the ends with metal plates [2]

This implies that $B_0^+ = -B_0^-$, $kd = p\pi$ and : $B_z = -2jB_0^+ \sin\left(\frac{p\pi z}{d}\right)\cos\left(\frac{m\pi x}{a}\right)\cos\left(\frac{n\pi y}{b}\right)e^{j\omega t}$



The frequency of a waveguide mode earlier was

$$ck = \sqrt{\omega^2 - \omega_{mn}^2}$$

Now, k is restricted, so

$$\omega_{0} = c\pi \sqrt{\left(\frac{p}{d}\right)^{2} + \left(\frac{m}{a}\right)^{2} + \left(\frac{n}{b}\right)^{2}}$$

Waveguide frequency

Losses in a resonator

• Perfect metal \rightarrow No losses

• Finite $\sigma \Rightarrow$ Boundary conditions get slightly modified \Rightarrow System supports small range of frequencies around ω_0

Waveguide resonator [2]



 $Q = \omega_0 \frac{Energy \, Stored}{Power \, lost}$



(b) Input spectrum



(c) Output spectrum

Waveguide resonator [2]



(b) Input spectrum



(c) Output spectrum

Calculating resonator loss allows Q calc.

Resonators in the real world

• Fiber-Bragg mirrors [3]





Resonators in the real world

• Fiber-based resonator [4]







The presence of a photon encourages another photon to be emitted: Stimulated emission

Only 2 level system,
$$\frac{N_2}{N_1} = \exp(-\frac{\Delta E}{kT})$$



So, the resonant frequency of the resonator must match the energy transition of the gain/laser medium!







Gain medium in the fiber? [6]

 Introduce a rare Earth ion into the fiber: Called "doping" – Erbium Doped Fiber



Applications [7]

- As a gain medium in a fiber-laser
- As an amplifier in fiber optics: backbone of ALL telecom!



VCSEL

• Vertical Cavity Surface Emitting Laser [8]



References

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- [2] Ulaby et al., Fundamentals of applied electromagnetics, 6th Ed.
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- [4] <u>http://spie.org/x8609.xml</u>
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