

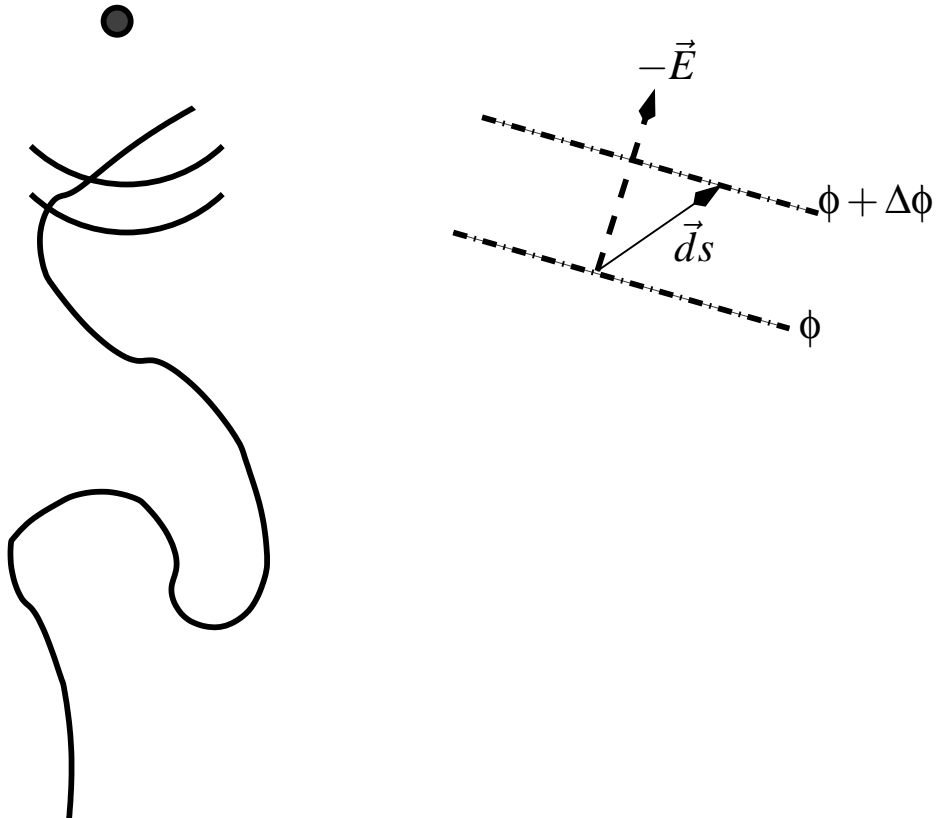
Work done to bring a charge from infinity

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Consider a charge Q at the origin. The field due to that charge is

$$\vec{E}(\vec{r}) = \frac{Q}{4\pi\epsilon_0} \frac{\hat{r}}{r^2}$$

Another charge q is brought along an arbitrary path to a point \vec{r} .



The work done to bring the charge to its final point is given by

$$\int_{\infty}^{\vec{r}} (-q\vec{E}) \cdot d\vec{s}' = \left(-\frac{qQ}{4\pi\epsilon_0} \right) \int_{\infty}^{\vec{r}} \frac{\hat{r}' \cdot d\vec{s}'}{r'^2} = \left(\frac{qQ}{4\pi\epsilon_0} \right) \int_r^{\infty} \frac{dr'}{r'^2}$$

The last integral is really over $r(\vec{s})$. However, the integrand is a perfect derivative, and the integral can be performed directly to yield

$$W = \left(\frac{qQ}{4\pi\epsilon_0} \right) \frac{1}{r}$$

Thus the result is not a function of the path that got us there.

If you look at the derivation, any force that had the form $\vec{F} = f(r)\hat{r}$ would have yielded the same independence of path. Such forces are called central forces and are important for this reason. From a vector calculus point of view, they are force fields that have zero curl (try it out). Thus, given any two points P and Q , we can apply Stoke's Theorem to a closed loop consisting of an arbitrary path from P to Q and another arbitrary path from Q to P .

$$\int_P^Q \vec{F} \cdot d\vec{S}_1 + \int_Q^P \vec{F} \cdot d\vec{S}_2 = \oint \vec{F} \cdot d\vec{S} = \int \int \nabla \times \vec{F} \cdot d\vec{S} = 0$$

which means that the integral is path independent:

$$\int_P^Q \vec{F} \cdot d\vec{S}_1 = \int_P^Q \vec{F} \cdot d\vec{S}_2$$

Thus, it is the central nature of Coloumb's Law that yields $\nabla \times \vec{E} = 0$ and permits us to derive \vec{E} from a scalar field ϕ .