

# The Wave Equation: Initial Value Example

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We start with the wave equation in 3 dimensions:

$$\nabla \left( \nabla \cdot \vec{E} \right) - \nabla^2 \vec{E} = -\omega^2 \mu \epsilon \vec{E}$$

which has the solution

$$\vec{E}(\vec{r}, t) = \iiint \left[ \vec{A}(\vec{k}) e^{j(kct - \vec{k} \cdot \vec{r})} + \vec{B}(\vec{k}) e^{j(-kct - \vec{k} \cdot \vec{r})} \right] d^3k \quad (1)$$

Suppose we are given that at  $t = 0$ ,

$$\begin{aligned} \vec{E} &= 2 \sin(k_0 z) \hat{x} + 3 \cos(k_1 y) \hat{z} \\ \eta \vec{H} &= \cos(k_1 y) \hat{x} - \sin(k_2 x) \hat{y} \end{aligned}$$

How do we find the fields at later time? At  $t = 0$  we have

$$\vec{E} = \iiint \left[ \vec{A}(\vec{k}) e^{j(-\vec{k} \cdot \vec{r})} + \vec{B}(\vec{k}) e^{j(-\vec{k} \cdot \vec{r})} \right] d^3k$$

and we have from Faraday's law that

$$\begin{aligned} \vec{H} &= -\frac{\nabla \times \vec{E}}{\pm j \omega \mu} \\ \vec{H} &= \iiint \left[ \frac{\vec{k} \times \vec{A}(\vec{k})}{\omega \mu} e^{j(-\vec{k} \cdot \vec{r})} - \frac{\vec{k} \times \vec{B}(\vec{k})}{\omega \mu} e^{j(-\vec{k} \cdot \vec{r})} \right] d^3k \end{aligned}$$

where the minus sign comes from the time derivative. We can also write this as

$$\eta \vec{H} = \iiint \left[ \hat{k} \times \vec{A} e^{j(-\vec{k} \cdot \vec{r})} - \hat{k} \times \vec{B} e^{j(-\vec{k} \cdot \vec{r})} \right] d^3k$$

If  $\vec{A}$  and  $\vec{B}$  had constant directions, clearly, we can write from Eq. 1

$$\vec{E}(\vec{r}, t) = \vec{f}(\vec{k} \cdot \vec{r} - kct) + \vec{g}(\vec{k} \cdot \vec{r} + kct)$$

Then, the magnetic field is given by

$$\vec{H}(\vec{r}, t) = \frac{1}{\eta} \left[ \hat{k} \times \vec{f}(\vec{k} \cdot \vec{r} - kct) - \hat{k} \times \vec{g}(\vec{k} \cdot \vec{r} + kct) \right]$$

Then,

$$\vec{E} \times \vec{H} = \frac{1}{\eta} \left( \vec{f} + \vec{g} \right) \times \left( \hat{k} \times \left( \vec{f} - \vec{g} \right) \right) = \frac{1}{\eta} \left[ \vec{f} \times \left( \hat{k} \times \vec{f} \right) - \vec{g} \times \left( \hat{k} \times \vec{g} \right) \right]$$

where we have ignored the cross terms assuming them to have zero mean. For a uniform medium, the fields are perpendicular to  $\vec{k}$ , and the result is

$$\begin{aligned}\vec{f} \times (\vec{k} \times \vec{f}) &= \epsilon_{ijl} \hat{x}_i f_j (\epsilon_{lmn} k_m f_n) \\ &= (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) \hat{x}_i f_j k_m f_n \\ &= \vec{k} |\vec{f}|^2 - \vec{f} (\vec{f} \cdot \vec{k}) \\ &= \vec{k} |\vec{f}|^2\end{aligned}$$

So,

$$\vec{E} \times \vec{H} = \frac{\hat{k}}{\eta} (|\vec{f}|^2 - |\vec{g}|^2)$$

as expected. The net flux of power along  $\vec{k}$  is the difference in the power flux in the forward going wave and the backward going wave.

Now consider the initial conditions. There are only three  $k$  values present, namely  $k_0 \hat{z}$ ,  $k_1 \hat{y}$  and  $k_2 \hat{x}$ . The directions are got from the multiplying coordinate, since it has to come out of  $\vec{k} \cdot \vec{r}$ . Since the equation is linear, we can break the problem into three problems, one in each  $k$  value.

$$k_0 \hat{z}: \vec{E} = 2 \sin(k_0 z) \hat{x}, \vec{H} = 0 \text{ (}\vec{H} \text{ must be along } \hat{z} \times \hat{x} = \hat{y}\text{)}$$

$$k_1 \hat{y}: \vec{E} = 3 \cos(k_1 y) \hat{z}, \eta \vec{H} = \cos(k_1 y) \hat{x} \text{ (The directions of } \vec{E} \text{ and } \vec{H} \text{ are consistent with each other)}$$

$$k_2 \hat{x}: \vec{E} = 0, \eta \vec{H} = -\sin(k_2 x) \hat{y} \text{ (}\vec{E} \text{ must be along } \hat{z}, \text{ so that } \hat{x} \times \hat{z} = -\hat{y}\text{)}$$

We now solve these problems in turn.

1. For  $k_0 \hat{z}$ : Since  $\vec{H}$  is zero at  $t = 0$ , we require  $A - B = 0$ , i.e.,  $A(k_0) = B(k_0)$  and  $A(-k_0) = B(-k_0)$ . By inspection we can write

$$\vec{E}(\vec{r}, t) = (\sin(k_0 z - \omega t) + \sin(k_0 z + \omega t)) \hat{x}$$

Then, the magnetic field is given by

$$\eta \vec{H}(\vec{r}, t) = (\sin(k_0 z - \omega t) - \sin(k_0 z + \omega t)) \hat{y}$$

2. For  $k_2 \hat{x}$ :  $\eta \vec{H} = -\sin(k_2 x) \hat{y}$ .  $\vec{E}$  must be along  $\hat{z}$ , and is zero at  $t = 0$ . Hence,  $\vec{A}(k_2) = -\vec{B}(k_2)$ . Again, by inspection, we can write

$$\eta \vec{H}(\vec{r}, t) = (-\sin(k_2 x - \omega t) - \sin(k_2 x + \omega t)) \frac{\hat{y}}{2}$$

with

$$\vec{E}(\vec{r}, t) = (\sin(k_2 x - \omega t) - \sin(k_2 x + \omega t)) \frac{\hat{z}}{2}$$

where I have used  $\hat{z} \times -\hat{y} = \hat{x}$  to obtain the direction of  $\vec{E}$ .

3. For  $k_1 \hat{y}$ : We have  $\hat{z} \times \hat{x} = \hat{y}$ , so the directions of  $\vec{E}$  and  $\vec{H}$  are consistent with each other. At  $t = 0$ ,

$$\begin{aligned}\vec{E} &= 3 \cos(k_1 y) \hat{z} \\ \vec{H} &= \cos(k_1 y) \hat{x}\end{aligned}$$

Thus, at  $k_1$ ,

$$\begin{aligned} A + B &= 3 \\ A - B &= 1 \end{aligned}$$

which yields  $A(k_1) = 2, B(k_1) = 1$ . Thus

$$\begin{aligned} \vec{E}(\vec{r}, t) &= [2 \cos(k_1 y - \omega t) + \cos(k_1 y + \omega t)] \hat{z} \\ \vec{H}(\vec{r}, t) &= [2 \cos(k_1 y - \omega t) - \cos(k_1 y + \omega t)] \hat{x} \end{aligned}$$

4. Suppose that in part 3,  $\vec{H}(\vec{r}, 0) = \cos(k_1 y) \hat{z}$ . What then? If the wave is to propagate along  $\vec{y}$ , the  $E_z$  requires a  $H_x$  and a  $H_z$  requires a  $-E_x$ . Thus the solution then have been

$$\begin{aligned} \vec{E}(\vec{r}, t) &= \frac{3}{2} [\cos(k_1 y - \omega t) + \cos(k_1 y + \omega t)] \hat{z} \\ &\quad - \frac{1}{2} [\cos(k_1 y - \omega t) - \cos(k_1 y + \omega t)] \hat{x} \\ \eta \vec{H}(\vec{r}, t) &= \frac{3}{2} [\cos(k_1 y - \omega t) - \cos(k_1 y + \omega t)] \hat{x} \\ &\quad + \frac{1}{2} [\cos(k_1 y - \omega t) + \cos(k_1 y + \omega t)] \hat{z} \end{aligned}$$

Thus the missing components of the fields would have appeared once  $t > 0$ .

5. Suppose that in part 3,  $\vec{H}(\vec{r}, 0) = \sin(k_1 y) \hat{x}$ . Then, we would have had

$$\begin{aligned} f(z) + g(z) &= 3 \cos(k_1 y) \\ f(z) - g(z) &= \sin(k_1 y) \end{aligned}$$

Clearly the solution is

$$\begin{aligned} f(z) &= \frac{3}{2} \cos(k_1 y) + \frac{1}{2} \sin(k_1 y) \\ g(z) &= \frac{3}{2} \cos(k_1 y) - \frac{1}{2} \sin(k_1 y) \end{aligned}$$

i.e.,

$$\begin{aligned} \vec{E}(\vec{r}, t) &= \frac{3}{2} [\cos(k_1 y - \omega t) + \cos(k_1 y + \omega t)] \hat{z} \\ &\quad + \frac{1}{2} [\sin(k_1 y - \omega t) - \sin(k_1 y + \omega t)] \hat{x} \\ \vec{H}(\vec{r}, t) &= \frac{3}{2} [\cos(k_1 y - \omega t) - \cos(k_1 y + \omega t)] \hat{x} \\ &\quad + \frac{1}{2} [\sin(k_1 y - \omega t) + \sin(k_1 y + \omega t)] \hat{z} \end{aligned}$$