

Can there be discontinuities in the phase response other than  $\pi$ ?

Recall that

$$H(e^{j\omega}) = |H(e^{j\omega})| e^{j\theta(\omega)} = H_R(e^{j\omega}) + j H_I(e^{j\omega})$$

$$|H(e^{j\omega})|^2 = H_R^2(e^{j\omega}) + H_I^2(e^{j\omega})$$

$$\theta(\omega) = \begin{cases} \tan^{-1} \frac{H_I(e^{j\omega})}{H_R(e^{j\omega})} & H(e^{j\omega}) \neq 0 \\ \text{undefined} & H(e^{j\omega}) = 0 \end{cases}$$

By definition,  $\theta(\omega) \in (-\pi, \pi]$  i.e.,  $-\pi < \theta(\omega) \leq \pi$

Note that  $e^{j\omega} \in \text{RoC}$  of  $H(z)$  and  $H(z)$  is a rational transfer function

Hence  $H_R(e^{j\omega})$  and  $H_I(e^{j\omega})$  are continuous functions of  $\omega$ .

Discontinuities in  $\theta(\omega)$  will occur in the following two cases:

(i) At points where  $H_I(e^{j\omega_0}) = 0$  and  $H_R(e^{j\omega_0}) < 0$ ,  $\theta(\omega_0) = \pi$ .

If  $H_I(e^{j\omega_0^-}) < 0$  or  $H_I(e^{j\omega_0^+}) < 0$  (or both), then

$\theta(\omega_0^-) = -\pi$ ,  $\theta(\omega_0^+) = -\pi$ . Hence, the phase jumps by  $2\pi$ ,

i.e.,  $\pi - (-\pi) = 2\pi$ .

(ii) If  $H(e^{j\omega_0}) = 0$ , then  $\theta(\omega_0)$  is undefined and hence phase cannot be continuous at that point.

The no. of points at which the phase can become discontinuous is finite because  $H(z)$  is rational.

Jumps of  $2\pi$  in  $\theta(\omega)$  can be removed by adding or subtracting integer multiples of  $2\pi$  - called **PHASE UNWRAPPING**

If we define  $\theta(\omega)$  suitably at points where  $H(e^{j\omega}) = 0$ , is it possible to get rid of discontinuities in  $\theta(\omega)$ ?

The answer is NO because we cannot get rid of jumps of  $\pi$  (odd multiples) by phase unwrapping.

Nevertheless, there is a way to make the phase continuous for systems with rational transfer functions.

Crossing a zero on the unit circle introduces a sign change. But  $|H(e^{j\omega})|$  is constrained to be non-negative. Hence the phase is forced to jump by  $\pi$ .

If we replace  $|H(e^{j\omega})|$  by  $A(\omega)$ , where  $A(\omega) \in \mathbb{R}$ , then

the change of sign can be absorbed in  $A(\omega)$  and the phase can remain continuous.

Hence, we decompose the frequency response as

$$H(e^{j\omega}) = \underbrace{A(\omega)}_{\text{real-valued, i.e., can take on both +ve and -ve values.}} e^{j\underbrace{\phi(\omega)}_{\text{continuous phase function}}}$$

The decomposition  $A(\omega)e^{j\phi(\omega)}$  is not unique because  $A(\omega)e^{j\phi(\omega)}$  is the same as  $-A(\omega)e^{j(\phi(\omega)+\pi)}$ .

The decomposition can be made *unique* if we enforce the following constraint:

$$0 \leq \phi(0) < \pi$$

### Example

Recall the example where  $h[n] = 1$   $0 \leq n \leq 2N$

$$H(e^{j\omega}) = e^{-j\omega N} \frac{\text{Sin}(2N+1)\omega/2}{\text{Sin}\omega/2}$$

The usual magnitude-phase decomposition resulted in a  $\theta(\omega)$  that had jumps of  $\pi$  at the zero crossings. If we replace  $|H(e^{j\omega})|$  by  $A(\omega)$  where

$$A(\omega) = \frac{\sin(2N+1)\omega/2}{\sin\omega/2}$$

then  $\phi(\omega) = -N\omega$ , which is now continuous. Note that  $A(\omega)$  given above now takes on both +ve and -ve values.

Phase response of a single complex pole is the negative of the phase response of a single complex zero.

Overall response is the result of the responses due to the individual poles and zeros.

Take a look at Example 5.10 in Oppenheim and Schaffer's, "Discrete-Time Signal Processing" (2nd edition)

See also MATLAB's `UNWRAP` and `ANGLE` commands.



When we were discussing causal signals, we saw that  $H_R(e^{j\omega})$  and  $H_I(e^{j\omega})$  are not independent but related. Does it mean  $|H(e^{j\omega})|$  and  $\theta(\omega)$  are also related?

Consider

$$H_I(z) = 1 - az^{-1} = 1 - re^{j\theta}z^{-1}$$

$$|H_I(e^{j\omega})|^2 = 1 + r^2 - 2r \cos(\omega - \theta)$$

$$\theta_I(\omega) = \tan^{-1} \frac{r \sin(\omega - \theta)}{1 - r \cos(\omega - \theta)}$$

Now let  $H_2(z) = -a^* + z^{-1} = -re^{-j\theta} + z^{-1}$

$$H_2(e^{j\omega}) = -re^{-j\theta} + e^{-j\omega}$$

$$H_2(e^{j\omega}) \cdot H_2^*(e^{j\omega}) = (e^{-j\omega} - re^{-j\theta})(e^{j\omega} - re^{j\theta})$$

$$= 1 + r^2 - 2r \cos(\omega - \theta)$$

$$= |H_1(e^{j\omega})|^2 \quad \text{same magnitude response!}$$

Phase response is different:

$$\theta_2(\omega) = \tan^{-1} \frac{r \sin \theta - \sin \omega}{\cos \omega - r \cos \theta}$$

Zero of  $H_1(z)$  is at  $re^{j\theta}$

zero of  $H_2(z)$  is at  $\frac{1}{r}e^{j\theta}$  i.e., the old zero is reflected about the unit circle.

Note:  $H_2(z) = -a^* + z^{-1}$

$$= -a^* \left[ 1 - \frac{1}{a^*} z^{-1} \right]$$

$$= \underbrace{-re^{-j\theta}} \left[ 1 - \frac{1}{r} e^{j\theta} z^{-1} \right]$$

scale factor needed to make the magnitude identical

