

Computational Electromagnetics : Finite Difference Time Domain Methods – Sources

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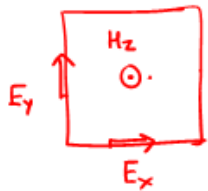
Topics in this module

- ① Current Sources
- ② Indirect Sources: Scattering problems
- ③ Summary of FDTD

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(2D) TE case $\rightarrow (E_x, E_y, H_z)$



$$\nabla \times \vec{H}^{n-1/2} = \epsilon \dot{\vec{E}}^{n-1/2} + \vec{J}^{n-1/2}$$

We need to know $\vec{J} = (J_x, J_y, J_z)$ as fn (r, t)

at time instances: $n-1/2$

at space instances: E field locations.

Easy to implement.

Let $\Delta x, \Delta t, \alpha$

Volume current excitation

$\vec{J}(r, t)$

Relation between current source and $\Delta t, \Delta x$?

$\vec{J}(t) \xleftrightarrow{\mathcal{F}} \vec{J}(f)$, say bandlimited.

$$\vec{J}(f) = 0, f > f_0$$

Nyquist thm: correctly represent $\vec{J}(t)$

$$\Delta t \leq \frac{1}{2f_0}$$

\Rightarrow High BW current source
 \Rightarrow space discretization fixed.

At the same time,

$$\text{Courant factor: } \alpha = \frac{c \Delta t}{\Delta x}$$

f_0 fixed $\Rightarrow \Delta t$ fixed $\Rightarrow \Delta x$ fixed.

meep

Other implementation issues

1) Gaussian current source.

$$g(t) = \exp\left(-\left(\frac{t-t_0}{t_w}\right)^2\right) \xleftrightarrow{\mathcal{F}} t_w \sqrt{\pi} \exp\left[-(\pi t_w f)^2\right] \exp[-j 2\pi f t_0]$$

What is $f_{bw} = \frac{1}{\pi t_w}$. To be safe $f_0 = 2 f_{bw} \Rightarrow f_0 = f_{max} = \frac{2}{\pi t_w}$
 \Rightarrow fixes Δx .

2) At start, $t=0$, $g(0) = e(-\frac{t_0}{t_w})^2$. Minimize high values of $g(0)$

Make t_0 large. eg. $\underline{t_0} \approx 4 t_w \Rightarrow$ longer simulations.

3) How long to run the sim? long enough e.g. $4 t_w \times 2$.

Common mistake e.g. $\underline{T} \approx 2 t_w$.

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No " $J(r, t)$ " term in scattering problems

$$\vec{E}_{\text{tot}} = \vec{E}_{\text{inc}} + \vec{E}_{\text{scat}}$$



1) In region I, \vec{E}_{scat} satisfies Maxwell's eqns.

\Rightarrow Apply FDTD update eqns to \vec{E}_s

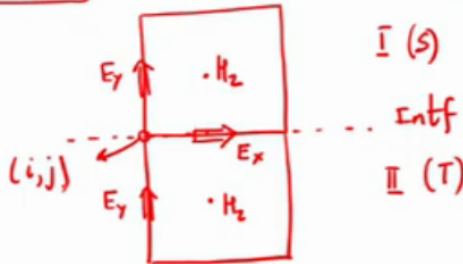
2) In region II, only \vec{E}_{tot} satisfies ME

\Rightarrow Discontinuity in variables at interface by E_{inc} .
 \hookrightarrow Method to incorporate E_{inc} (source)

TE pol.

$$1) \quad \epsilon \frac{\partial E_y}{\partial t} = -\frac{\partial H_z}{\partial x}$$

$$2) \quad \epsilon \frac{\partial E_x}{\partial t} = \frac{\partial H_z}{\partial y}$$



Yee cell \int

$$3) \quad \frac{\partial H_z}{\partial t} = \frac{\partial E_x}{\partial y} - \frac{\partial E_y}{\partial x}$$



In S cell.

Fix via total/scattered formulation

$$\text{Eqn (1)} \Rightarrow \epsilon \left(\frac{E_{y,s}^n - E_{y,s}^{n-1}}{\Delta t} \right) = \frac{H_{z,s}^{n-1/2}(i-1/2, j+1/2) - H_{z,s}^{n-1/2}(i+1/2, j+1/2)}{\Delta x} \quad \text{unchanged.}$$

$$\begin{aligned} \text{Eqn (2)} \Rightarrow \epsilon \left(\frac{E_{x,t}^n - E_{x,t}^{n-1}}{\Delta t} \right) &= \frac{H_{z,t}(i+1/2, j+1/2) - H_{z,t}(i+1/2, j-1/2)}{\Delta y} \\ &= \frac{H_{z,s}(i+1/2, j+1/2) - H_{z,t}(i+1/2, j-1/2)}{\Delta y} + \frac{H_{z,i}(i+1/2, j+1/2)}{\Delta y} \end{aligned}$$

But we are storing $H_{z,s}$
& not $H_{z,t}$ in $\underline{\Sigma}$.

interpret as "current" term

② Choose for $E_x \rightarrow$ scat or total
on interface.

① \therefore incident field gets introduced at
the interface



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flow of FDTD.

Summary of FDTD

- ↳ Maxwell's eqns: finite differences (Taylor's thm)
- ↳ Stencil (Yee cell) → update eqns
- ↳ Analysis - Convergence, stability (Courant)
 - Accuracy, dispersion (numerical)
- ↳ Materials - simple (Ohm's law), dispersive media (Debye)
 $D = \epsilon E$ (freq domain)
- ↳ Absorbing boundary conds
- ↳ PMLs (stretched coordinates)
- ↳ Sources

MEEP ..



Topics that were covered in this module

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References:

- * Ch 12 of Computational Methods for Electromagnetics - Peterson, Ray, Mitra
- * Computational Electrodynamics: The Finite-Difference Time-Domain Method – Allen Taflove (the 'Bible' for FDTD)